# The Inverse Problem for Periodic Travelling Waves of the Linear 1D Shallow-Water Equations

### Robert Hakl

Institute of Mathematics CAS, Czech Republic

Czech-Georgian Workshop Brno, July 2-4, 2024

R. Hakl

Czech-Georgian Workshop Brno, July 2-4, 2024

1/17

The motion of small amplitude waves of a water layer with variable depth along the x-axis is described by the equations of the shallow water theory

$$rac{\partial\eta}{\partial t}+rac{\partial}{\partial x}\left[h(x)u
ight]=0,\qquad rac{\partial u}{\partial t}+grac{\partial\eta}{\partial x}=0, \tag{1}$$

Czech-Georgian Workshop Brno, July 2-4, 2024

2/17

The motion of small amplitude waves of a water layer with variable depth along the x-axis is described by the equations of the shallow water theory

$$rac{\partial\eta}{\partial t}+rac{\partial}{\partial x}\left[h\left(x
ight)u
ight]=0,\qquad rac{\partial u}{\partial t}+grac{\partial\eta}{\partial x}=0, \tag{1}$$

Czech-Georgian Workshop Brno, July 2-4, 2024

2/17

where

•  $\eta(x,t)$  is the vertical water surface elevation

The motion of small amplitude waves of a water layer with variable depth along the x-axis is described by the equations of the shallow water theory

$$rac{\partial\eta}{\partial t}+rac{\partial}{\partial x}\left[h\left(x
ight)u
ight]=0,\qquad rac{\partial u}{\partial t}+grac{\partial\eta}{\partial x}=0, \tag{1}$$

Czech-Georgian Workshop Brno, July 2-4, 2024

2/17

where

- $\eta(x,t)$  is the vertical water surface elevation
- u(x,t) is the depth-averaged water flow velocity (also called wave velocity)

The motion of small amplitude waves of a water layer with variable depth along the x-axis is described by the equations of the shallow water theory

$$rac{\partial\eta}{\partial t}+rac{\partial}{\partial x}\left[h\left(x
ight)u
ight]=0,\qquad rac{\partial u}{\partial t}+grac{\partial\eta}{\partial x}=0, \tag{1}$$

Czech-Georgian Workshop Brno, July 2-4, 2024

2/17

where

- $\eta(x,t)$  is the vertical water surface elevation
- u(x,t) is the depth-averaged water flow velocity (also called wave velocity)
- h(x) is the unperturbed water depth

The motion of small amplitude waves of a water layer with variable depth along the x-axis is described by the equations of the shallow water theory

$$rac{\partial\eta}{\partial t}+rac{\partial}{\partial x}\left[h\left(x
ight)u
ight]=0,\qquad rac{\partial u}{\partial t}+grac{\partial\eta}{\partial x}=0, \tag{1}$$

Czech-Georgian Workshop Brno, July 2-4, 2024

2/17

where

- $\eta(x,t)$  is the vertical water surface elevation
- u(x,t) is the depth-averaged water flow velocity (also called wave velocity)
- h(x) is the unperturbed water depth
- g is the gravity acceleration (we assume without loss of generality that g = 1)

The motion of small amplitude waves of a water layer with variable depth along the x-axis is described by the equations of the shallow water theory

$$rac{\partial\eta}{\partial t}+rac{\partial}{\partial x}\left[h\left(x
ight)u
ight]=0,\qquad rac{\partial u}{\partial t}+grac{\partial\eta}{\partial x}=0, \tag{1}$$

where

- $\eta(x,t)$  is the vertical water surface elevation
- u(x,t) is the depth-averaged water flow velocity (also called wave velocity)
- h(x) is the unperturbed water depth
- g is the gravity acceleration (we assume without loss of generality that g = 1)

The shallow water equations conform a system of coupled PDEs of first order that can be easily decoupled into a single wave equation for the water flow velocity

$$rac{\partial^2 u}{\partial t^2} - rac{\partial^2}{\partial x^2} \left[ h(x) u 
ight] = 0,$$
 (2)

or for the surface elevation

$$\frac{\partial^2 \eta}{\partial t^2} - \frac{\partial}{\partial x} \left[ h(x) \frac{\partial \eta}{\partial x} \right] = 0.$$
(3)

R. Hakl

2/17

The water flow velocity equation:

$$rac{\partial^2 u}{\partial t^2} - rac{\partial^2}{\partial x^2} \left[h(x)u
ight] = 0,$$
 (2)

The surface elevation equation:

$$\frac{\partial^2 \eta}{\partial t^2} - \frac{\partial}{\partial x} \left[ h(x) \frac{\partial \eta}{\partial x} \right] = 0.$$
(3)

3/17

The water flow velocity equation:

$$\frac{\partial^2 u}{\partial t^2} - \frac{\partial^2}{\partial x^2} \left[ h(x) u \right] = 0, \qquad (2)$$

The surface elevation equation:

$$\frac{\partial^2 \eta}{\partial t^2} - \frac{\partial}{\partial x} \left[ h(x) \frac{\partial \eta}{\partial x} \right] = 0.$$
(3)

3/17

Czech-Georgian Workshop Brno, July 2-4, 2024

A travelling wave is a special solution of the form

 $q(x)\exp{\left(i\left[\omega t-\Psi(x)
ight]
ight)},$ 

where both q and  $\Psi$  are scalar real-valued functions.

The water flow velocity equation:

$$\frac{\partial^2 u}{\partial t^2} - \frac{\partial^2}{\partial x^2} \left[ h(x) u \right] = 0, \qquad (2)$$

The surface elevation equation:

$$\frac{\partial^2 \eta}{\partial t^2} - \frac{\partial}{\partial x} \left[ h(x) \frac{\partial \eta}{\partial x} \right] = 0.$$
(3)

3/17

Czech-Georgian Workshop Brno, July 2-4, 2024

A travelling wave is a special solution of the form

$$q(x)\exp{(i[\omega t-\Psi(x)])},$$

where both q and  $\Psi$  are scalar real-valued functions. In the related literature,

- q(x) is known as the amplitude of the travelling wave
- $\omega$  is the frequency
- $\Psi(x)$  is the phase

R. Hakl

The water flow velocity equation:

$$rac{\partial^2 u}{\partial t^2} - rac{\partial^2}{\partial x^2} \left[h(x)u\right] = 0,$$
 (2)

The surface elevation equation:

$$\frac{\partial^2 \eta}{\partial t^2} - \frac{\partial}{\partial x} \left[ h(x) \frac{\partial \eta}{\partial x} \right] = 0.$$
(3)

A travelling wave is a special solution of the form

 $q(x)\exp{(i[\omega t-\Psi(x)])},$ 

where both q and  $\Psi$  are scalar real-valued functions. In the related literature,

- q(x) is known as the amplitude of the travelling wave
- $\omega$  is the frequency
- $\Psi(x)$  is the phase

#### The inverse problem:

Given a prescribed amplitude q(x), can we determine a suitable bottom profile h(x) allowing the equation to admit a travelling wave with amplitude q(x)?

R. Hakl

3/17

 $C_T^+$  will denote the space of continuous scalar T-periodic functions with positive values.



 $C_T^+$  will denote the space of continuous scalar *T*-periodic functions with positive values. Given a fixed  $q \in C_T^+$ , we wonder if there exists  $h \in C_T^+$  such that Eq.

$$rac{\partial^2 u}{\partial t^2} - rac{\partial^2}{\partial x^2} \left[ h(x) u 
ight] = 0,$$
 (2)

4/17

 $C_T^+$  will denote the space of continuous scalar *T*-periodic functions with positive values. Given a fixed  $q \in C_T^+$ , we wonder if there exists  $h \in C_T^+$  such that Eq.

$$rac{\partial^2 u}{\partial t^2} - rac{\partial^2}{\partial x^2} \left[ h(x) u 
ight] = 0,$$
 (2)

4/17

Czech-Georgian Workshop Brno, July 2-4, 2024

has a travelling wave

$$u(x,t) = q(x) \exp\left(i\left[\omega t - \Psi(x)
ight]
ight).$$

 $C_T^+$  will denote the space of continuous scalar *T*-periodic functions with positive values. Given a fixed  $q \in C_T^+$ , we wonder if there exists  $h \in C_T^+$  such that Eq.

$$rac{\partial^2 u}{\partial t^2} - rac{\partial^2}{\partial x^2} \left[ h(x) u 
ight] = 0,$$
 (2)

has a travelling wave

$$u(x,t) = q(x) \exp \left( i \left[ \omega t - \Psi(x) \right] \right).$$
  
 $(hq)'' + \omega^2 q - hq \Psi'^2 = 0,$  (4)

$$2(hq)'\Psi' + hq\Psi'' = 0.$$
 (5)

4/17

 $C_T^+$  will denote the space of continuous scalar *T*-periodic functions with positive values. Given a fixed  $q \in C_T^+$ , we wonder if there exists  $h \in C_T^+$  such that Eq.

$$rac{\partial^2 u}{\partial t^2} - rac{\partial^2}{\partial x^2} \left[h(x)u\right] = 0,$$
 (2)

has a travelling wave

$$u(x,t)=q(x)\expig(i\left[\omega t-\Psi(x)
ight]ig).$$

$$(hq)'' + \omega^2 q - hq \Psi'^2 = 0, \tag{4}$$

$$2(hq)'\Psi' + hq\Psi'' = 0.$$
 (5)

4/17

Czech-Georgian Workshop Brno, July 2-4, 2024

From (5), we deduce that  $\left[(hq)^2\Psi'\right]'=0$ , and  $(hq)^2\Psi'$  is a conserved quantity.

 $C_T^+$  will denote the space of continuous scalar *T*-periodic functions with positive values. Given a fixed  $q \in C_T^+$ , we wonder if there exists  $h \in C_T^+$  such that Eq.

$$rac{\partial^2 u}{\partial t^2} - rac{\partial^2}{\partial x^2} \left[ h(x) u 
ight] = 0,$$
 (2)

has a travelling wave

$$u(x,t) = q(x) \exp\left(i \left[\omega t - \Psi(x)\right]\right).$$
  
 $(hq)'' + \omega^2 q - hq {\Psi'}^2 = 0,$  (4)

$$2(hq)'\Psi' + hq\Psi'' = 0.$$
 (5)

4/17

Czech-Georgian Workshop Brno, July 2-4, 2024

From (5), we deduce that  $\left[(hq)^2\Psi'\right]'=0$ , and  $(hq)^2\Psi'$  is a conserved quantity. This means that there exists  $\alpha \in \mathbb{R}$  such that

$$\left[h(x)q(x)\right]^{2}\Psi'(x)=lpha, \qquad \forall x\in\mathbb{R}.$$
 (6)

 $C_T^+$  will denote the space of continuous scalar *T*-periodic functions with positive values. Given a fixed  $q \in C_T^+$ , we wonder if there exists  $h \in C_T^+$  such that Eq.

$$rac{\partial^2 u}{\partial t^2} - rac{\partial^2}{\partial x^2} \left[ h(x) u 
ight] = 0,$$
 (2)

has a travelling wave

$$u(x,t) = q(x) \exp\left(i \left[\omega t - \Psi(x)\right]\right).$$
  
 $(hq)'' + \omega^2 q - hq {\Psi'}^2 = 0,$  (4)

$$2(hq)'\Psi' + hq\Psi'' = 0.$$
 (5)

From (5), we deduce that  $\left[(hq)^2\Psi'\right]'=0$ , and  $(hq)^2\Psi'$  is a conserved quantity. This means that there exists  $\alpha\in\mathbb{R}$  such that

$$\left[h\left(x
ight)q\left(x
ight)
ight]^{2}\Psi'(x)=lpha,\qquadorall x\in\mathbb{R}.$$
 (6)

Now, we insert (6) into (4) and arrive to a single second order ODE

$$(hq)'' + \omega^2 q = \frac{\alpha^2}{(hq)^3}.$$
 (7)

4/17

$$(hq)'' + \omega^2 q = \frac{\alpha^2}{(hq)^3} \tag{7}$$

### Theorem 1

There exists a solution  $h \in C_T^+$  of (7) for any  $\alpha \neq 0$ ,  $\omega \neq 0$ .





$$(hq)'' + \omega^2 q = \frac{\alpha^2}{(hq)^3} \tag{7}$$

5/17

### Theorem 1

There exists a solution  $h \in C_T^+$  of (7) for any  $\alpha \neq 0$ ,  $\omega \neq 0$ .

**Proof.** By introducing the change of variables y = hq into (7), we get the equation

$$y^{\prime\prime}+\omega^2 q=rac{lpha^2}{y^3}.$$

Now, the result is a direct consequence of Theorem 3.12 in

[LS] A.C. Lazer, S. Solimini, On periodic solutions of nonlinear differential equations with singularities, Proc. American Math. Society 99, No. 1, 1987.

$$\frac{\partial^2 \eta}{\partial t^2} - \frac{\partial}{\partial x} \left[ h(x) \frac{\partial \eta}{\partial x} \right] = 0 \tag{3}$$



$$\frac{\partial^2 \eta}{\partial t^2} - \frac{\partial}{\partial x} \left[ h(x) \frac{\partial \eta}{\partial x} \right] = 0 \tag{3}$$

Czech-Georgian Workshop Brno, July 2-4, 2024

6/17

has a travelling wave of the form

$$\eta(x,t) = q(x) \exp ig( i \left[ \omega t - \Psi(x) 
ight] ig).$$

$$\frac{\partial^2 \eta}{\partial t^2} - \frac{\partial}{\partial x} \left[ h(x) \frac{\partial \eta}{\partial x} \right] = 0 \tag{3}$$

has a travelling wave of the form

$$\eta(x,t) = q(x) \exp\left(i \left[\omega t - \Psi(x)\right]\right).$$

$$(hq')' + \omega^2 q - hq \Psi'^2 = 0,$$
(8)

$$(hq\Psi')' + hq'\Psi' = 0. (9)$$

$$\frac{\partial^2 \eta}{\partial t^2} - \frac{\partial}{\partial x} \left[ h(x) \frac{\partial \eta}{\partial x} \right] = 0 \tag{3}$$

has a travelling wave of the form

$$\eta(x,t)=q(x)\expig(i\left[\omega t-\Psi(x)
ight]ig).$$

$$(hq')' + \omega^2 q - hq \Psi'^2 = 0, \qquad (8)$$

$$(hq\Psi')' + hq'\Psi' = 0.$$
 (9)

6/17

Czech-Georgian Workshop Brno, July 2-4, 2024

Now, the conserved quantity coming from (9) is

$$h(x)q(x)^2\Psi'(x)=lpha,\qquad orall x\in\mathbb{R}.$$

$$\frac{\partial^2 \eta}{\partial t^2} - \frac{\partial}{\partial x} \left[ h(x) \frac{\partial \eta}{\partial x} \right] = 0 \tag{3}$$

has a travelling wave of the form

$$\eta(x,t)=q(x)\expig(i\left[\omega t-\Psi(x)
ight]ig).$$

$$(hq')' + \omega^2 q - hq \Psi'^2 = 0, \qquad (8)$$

$$(hq\Psi')' + hq'\Psi' = 0. (9)$$

Now, the conserved quantity coming from (9) is

$$h(x)q(x)^2\Psi'(x)=lpha,\qquad orall x\in\mathbb{R}.$$

Using this information in (8), we arrive at the equation

$$(hq')' + \omega^2 q = \frac{\alpha^2}{hq^3}.$$
 (10)

We assume that q(x) is a *T*-periodic and positive function of class  $C^2$  with a finite number of critical points in [0, T], all of them non-degenerate, that is, if q'(x) = 0 then  $q''(x) \neq 0$ . Under this assumption, we can divide the interval [0, T] into subintervals [a, b] such that q'(x) is of a constant sign on (a, b) and q'(a) = q'(b) = 0. Then, the substitution

$$u(x)=rac{(h(x)q'(x))^2}{2\omega^4} \qquad ext{for} \ x\in(a,b)$$

transforms

$$(hq')' + \omega^2 q = \frac{\alpha^2}{hq^3} \tag{10}$$

7/17

Czech-Georgian Workshop Brno, July 2-4, 2024

into the equation

$$u'(x) = rac{\lambda^2 q'(x)}{q^3(x)} - q(x) \operatorname{sgn}(q') \sqrt{2u(x)}$$
 for  $x \in (a, b),$  (12)

where  $\lambda = \alpha / \omega^2$ .

We consider an interval [a,b] such that  $q\in C^{2}\left([a,b];\mathbb{R}
ight)$  satisfies

 $q(x)>0 \quad ext{for} \ x\in [a,b], \quad q'(a)=0, \quad q'(b)=0, \quad q'(x)>0 \quad ext{for} \ x\in (a,b).$ 

and

$$q''(a) > 0, \qquad q''(b) < 0.$$
 (14)

8/17

We consider an interval [a,b] such that  $q\in C^2ig([a,b];\mathbb{R}ig)$  satisfies

q(x) > 0 for  $x \in [a, b]$ , q'(a) = 0, q'(b) = 0, q'(x) > 0 for  $x \in (a, b)$ . (13)

and

$$q''(a) > 0, \qquad q''(b) < 0.$$
 (14)

In such an interval, eq. (12) reads as

$$u'(x) = rac{\lambda^2 q'(x)}{q^3(x)} - q(x) \sqrt{2u(x)} \qquad ext{for } x \in (a,b).$$
 (15)

We consider an interval [a,b] such that  $q\in C^2ig([a,b];\mathbb{R}ig)$  satisfies

q(x) > 0 for  $x \in [a, b], q'(a) = 0, q'(b) = 0, q'(x) > 0$  for  $x \in (a, b).$  (13)

and

$$q''(a) > 0, \qquad q''(b) < 0.$$
 (14)

Czech-Georgian Workshop Brno, July 2-4, 2024

8/17

In such an interval, eq. (12) reads as

$$u'(x) = rac{\lambda^2 q'(x)}{q^3(x)} - q(x) \sqrt{2u(x)} \qquad ext{for } x \in (a,b).$$
 (15)

For technical reasons, we are going to embed this equation into

$$u'(x) = rac{\lambda^2 q'(x)}{q^3(x)} - q(x) \sqrt{2|u(x)|} \operatorname{sgn} u(x) \qquad ext{for } x \in [a,b].$$

Obviously, non-negative solutions of (15) and (16) are the same.

We consider an interval [a,b] such that  $q\in C^2ig([a,b];\mathbb{R}ig)$  satisfies

q(x) > 0 for  $x \in [a, b]$ , q'(a) = 0, q'(b) = 0, q'(x) > 0 for  $x \in (a, b)$ . (13)

and

$$q''(a) > 0, \qquad q''(b) < 0.$$
 (14)

In such an interval, eq. (12) reads as

$$u'(x) = rac{\lambda^2 q'(x)}{q^3(x)} - q(x) \sqrt{2u(x)} \qquad ext{for } x \in (a,b).$$
 (15)

For technical reasons, we are going to embed this equation into

$$u'(x) = rac{\lambda^2 q'(x)}{q^3(x)} - q(x) \sqrt{2|u(x)|} \operatorname{sgn} u(x) \qquad ext{for } x \in [a,b].$$
 (16)

Obviously, non-negative solutions of (15) and (16) are the same.

A solution to (16) is understood in the classical sense, that is, a function  $u \in C^1([a, b]; \mathbb{R})$  satisfying (16) for every  $x \in [a, b]$ . We will investigate the properties of a solution to (16) subject to the initial condition

$$u(a) = 0. \tag{17}$$

$$u'(x) = \frac{\lambda^2 q'(x)}{q^3(x)} - q(x)\sqrt{2|u(x)|} \operatorname{sgn} u(x) \quad \text{for } x \in [a, b],$$
(16)  
$$u(a) = 0. \quad (17)$$

$$u'(x) = rac{\lambda^2 q'(x)}{q^3(x)} - q(x)\sqrt{2|u(x)|}\operatorname{sgn} u(x)$$
 for  $x \in [a, b],$  (16)  
 $u(a) = 0.$  (17)

9/17

### QUESTIONS:

• When u(b) = 0?



$$u'(x) = \frac{\lambda^2 q'(x)}{q^3(x)} - q(x)\sqrt{2|u(x)|} \operatorname{sgn} u(x) \quad \text{for } x \in [a, b],$$
(16)  
$$u(a) = 0. \quad (17)$$

9/17

### QUESTIONS:

- When u(b) = 0?
- Do there exist one-sided limits

$$\ell_a \stackrel{def}{=} \lim_{x \to a+} \frac{\sqrt{2u(x)}}{q'(x)}, \qquad \ell_b \stackrel{def}{=} \lim_{x \to b-} \frac{\sqrt{2u(x)}}{q'(x)}$$
?

$$u'(x) = \frac{\lambda^2 q'(x)}{q^3(x)} - q(x)\sqrt{2|u(x)|} \operatorname{sgn} u(x) \quad \text{for } x \in [a, b],$$
(16)  
$$u(a) = 0. \quad (17)$$

9/17

### QUESTIONS:

- When u(b) = 0?
- Do there exist one-sided limits

$$\ell_a \stackrel{def}{=} \lim_{x \to a+} \frac{\sqrt{2u(x)}}{q'(x)}, \qquad \ell_b \stackrel{def}{=} \lim_{x \to b-} \frac{\sqrt{2u(x)}}{q'(x)}$$
?

• What the values  $\ell_a$  and  $\ell_b$  are equal to?

$$u'(x) = rac{\lambda^2 q'(x)}{q^3(x)} - q(x)\sqrt{2|u(x)|} \operatorname{sgn} u(x) \quad ext{ for } x \in [a,b], \ u(a) = 0.$$
 (16)

$$u'(x) = rac{\lambda^2 q'(x)}{q^3(x)} - q(x)\sqrt{2|u(x)|} \operatorname{sgn} u(x) \quad ext{ for } x \in [a, b],$$
 (16)  
 $u(a) = 0.$  (17)

#### Theorem 2

There exists a threshold  $\lambda_0 > 0$  such that

(i) if  $0 < |\lambda| < \lambda_0$ , the unique solution u of (16), (17) verifies u(b) = 0. Moreover,  $\ell_a$  and  $\ell_b$  are respectively the unique positive root of

$$y^{2} + \frac{q(a)}{q''(a)}y - \frac{\lambda^{2}}{q^{3}(a)q''(a)} = 0.$$
 (18)

and the smaller root of

$$y^{2} - \frac{q(b)}{|q''(b)|}y + \frac{\lambda^{2}}{q^{3}(b)|q''(b)|} = 0.$$
(19)

(ii) if |λ| = λ<sub>0</sub>, the unique solution u of (16), (17) verifies u(b) = 0. Moreover, ℓ<sub>a</sub> and ℓ<sub>b</sub> are respectively the unique positive root of (18) and a root of (19).
(iii) if |λ| > λ<sub>0</sub>, the unique solution u of (16), (17) verifies u(b) > 0.

Consider now an interval  $[ ilde{a}, ilde{b}]$  such that  $q\in C^2\left([ ilde{a}, ilde{b}];\mathbb{R}
ight)$  satisfies

 $q({m x})>0 \quad ext{for} \,\, {m x}\in [ ilde{a}, ilde{b}], \quad q'( ilde{a})=0, \quad q'( ilde{b})=0, \quad q'({m x})<0 \quad ext{for} \,\, {m x}\in ( ilde{a}, ilde{b}).$ 

and

$$q^{\prime\prime}( ilde{a}) < 0, \qquad q^{\prime\prime}( ilde{b}) > 0.$$

In this interval, Eq. (12) reads

$$u'(x)=rac{\lambda^2 q'(x)}{q^3(x)}+q(x)\sqrt{2u(x)} \qquad ext{for } x\in ( ilde{a}, ilde{b}).$$

Note that the function  $\tilde{q}(x) = q(-x)$  verifies (13) and (14) in the interval  $[a, b] = [-\tilde{b}, -\tilde{a}]$ . Moreover, v(x) = u(-x) satisfies the equation

$$v'(x)=rac{\lambda^2 ilde q'(x)}{ ilde q^3(x)}- ilde q(x)\sqrt{2v(x)} \qquad ext{for} \ x\in(a,b),$$

which is just (15). In conclusion, the case of an interval where q is decreasing can be reduced to the case studied above.

11 / 17

$$(hq')' + \omega^2 q = \frac{\alpha^2}{hq^3} \tag{10}$$

### Main Result

Let us assume that q is a T-periodic and positive function of class  $C^2$  with a finite number of critical points in [0, T], all of them non-degenerate, that is, if q'(x) = 0 then  $q''(x) \neq 0$ . Then, there exists a threshold  $\lambda_0 > 0$  such that

(i) there exists a positive T-periodic solution h of (10) provided  $0 < \left|\frac{\alpha}{\omega^2}\right| < \lambda_0$ ,

(ii) no positive T-periodic solution of (10) exists provided  $\left|\frac{\alpha}{\omega^2}\right| > \lambda_0$ .

Moreover,

$$\frac{q_*^5}{4|q_0|} < \lambda_0^2 \le \min\left\{\frac{q^5(b)}{4|q''(b)|} : q'(b) = 0, q''(b) < 0\right\},\tag{22}$$

where

$$q_* \stackrel{def}{=} \min\{q(x) : x \in [0,T]\}, \quad q_0 \stackrel{def}{=} \min\{q''(x) : x \in [0,T]\}.$$
 (23)

### Theorem 3

Let there exist positive constants  $\lambda_1$  and  $\lambda_2$  such that  $\lambda_1 \leq \lambda_2$ , and let  $v, w \in AC([a, b]; \mathbb{R})$  satisfy

$$egin{aligned} &v'(x) \geq rac{\lambda_1^2 q'(x)}{q^3(x)} - q(x) \sqrt{2|v(x)|} \, ext{sgn} \, v(x) & ext{ for a. e. } x \in [a,b], \ &w'(x) \leq rac{\lambda_2^2 q'(x)}{q^3(x)} - q(x) \sqrt{2|w(x)|} \, ext{sgn} \, w(x) & ext{ for a. e. } x \in [a,b], \ &v(a) \geq 0 \geq w(a), &v(b) = 0 = w(b) \ &\lim_{x o b^-} rac{\sqrt{2|w(x)|} \, ext{sgn} \, w(x)}{q'(x)} > y_1(\lambda_2), \end{aligned}$$

where  $y_1(\lambda_2)$  is the smaller root of (19) with  $\lambda = \lambda_2$ . Then, the threshold  $\lambda_0$  admits the estimate

$$\lambda_1 \le \lambda_0 \le \lambda_2. \tag{24}$$

R. Hakl

13 / 17

### Corollary

Let there exist positive constants  $\lambda_1$  and  $\lambda_2$  such that  $\lambda_1 \leq \lambda_2$ , let  $q \in C^2([a, b]; \mathbb{R})$ , and let  $\ell_1, \ell_2 \in C^1([a, b]; \mathbb{R})$  satisfy

$$\ell_i(x)>0 \qquad ext{for } x\in [a,b] \quad (i=1,2),$$

$$\ell_1(x) \Big( \ell_1'(x) q'(x) + \ell_1(x) q''(x) \Big) \geq rac{\lambda_1^2}{q^3(x)} - q(x) \ell_1(x) \qquad ext{for } x \in [a,b],$$

$$\ell_2(x) \Big( \ell_2'(x) q'(x) + \ell_2(x) q''(x) \Big) \le rac{\lambda_2^2}{q^3(x)} - q(x) \ell_2(x) \qquad ext{for } x \in [a, b],$$
 (27)  
 $\ell_2(b) > y_1(\lambda_2).$  (28)

where  $y_1(\lambda_2)$  is the smaller root of (19) with  $\lambda = \lambda_2$ . Then, the threshold  $\lambda_0$  admits the estimate

$$\lambda_1 \leq \lambda_0 \leq \lambda_2. \tag{24}$$

Czech-Georgian Workshop Brno, July 2-4, 2024
 Czech-Georgian Workshop Brno, July 2-4, 2024
 Vertex 14 / 17

Consider  $q(x) = 2 - \cos x$  for  $x \in [0, 2\pi]$ .

Consider  $q(x) = 2 - \cos x$  for  $x \in [0, 2\pi]$ . Then local extremes of q divide the interval  $[0, 2\pi]$  into two subintervals, in particular, we set  $T = 2\pi$ ,  $x_1 = 0$ ,  $x_2 = \pi$ ,  $x_1 + T = 2\pi$ .



Consider  $q(x) = 2 - \cos x$  for  $x \in [0, 2\pi]$ . Then local extremes of q divide the interval  $[0, 2\pi]$  into two subintervals, in particular, we set  $T = 2\pi$ ,  $x_1 = 0$ ,  $x_2 = \pi$ ,  $x_1 + T = 2\pi$ . Then we have

$$q'(x)>0 \quad ext{for} \ x\in(0,\pi), \qquad q'(x)<0 \quad ext{for} \ x\in(\pi,2\pi), \ q'(0)=q'(\pi)=q'(2\pi)=0, \qquad q''(0)=q''(2\pi)=1, \qquad q''(\pi)=-1.$$



Consider  $q(x) = 2 - \cos x$  for  $x \in [0, 2\pi]$ . Then local extremes of q divide the interval  $[0, 2\pi]$  into two subintervals, in particular, we set  $T = 2\pi$ ,  $x_1 = 0$ ,  $x_2 = \pi$ ,  $x_1 + T = 2\pi$ . Then we have

$$q'(x)>0 \quad ext{for } x\in(0,\pi), \qquad q'(x)<0 \quad ext{for } x\in(\pi,2\pi), \ q'(0)=q'(\pi)=q'(2\pi)=0, \qquad q''(0)=q''(2\pi)=1, \qquad q''(\pi)=-1.$$

Moreover, since q is symmetric with respect to  $\pi$ , we can easily conclude that the thresholds corresponding to each subinterval has the same value, i.e.,  $\lambda_0 = \lambda_{01} = \lambda_{02}$ .



Consider  $q(x) = 2 - \cos x$  for  $x \in [0, 2\pi]$ . Then local extremes of q divide the interval  $[0, 2\pi]$  into two subintervals, in particular, we set  $T = 2\pi$ ,  $x_1 = 0$ ,  $x_2 = \pi$ ,  $x_1 + T = 2\pi$ . Then we have

$$q'(x)>0 \quad ext{for } x\in(0,\pi), \qquad q'(x)<0 \quad ext{for } x\in(\pi,2\pi), \ q'(0)=q'(\pi)=q'(2\pi)=0, \qquad q''(0)=q''(2\pi)=1, \qquad q''(\pi)=-1.$$

Moreover, since q is symmetric with respect to  $\pi$ , we can easily conclude that the thresholds corresponding to each subinterval has the same value, i.e.,  $\lambda_0 = \lambda_{01} = \lambda_{02}$ . Thus, according to Main Result, the threshold  $\lambda_0$  satisfies the inequalities

$$0.25 = rac{1}{4} < \lambda_0^2 \leq rac{243}{4} = 60.75.$$

Consider  $q(x) = 2 - \cos x$  for  $x \in [0, 2\pi]$ . Then local extremes of q divide the interval  $[0, 2\pi]$  into two subintervals, in particular, we set  $T = 2\pi$ ,  $x_1 = 0$ ,  $x_2 = \pi$ ,  $x_1 + T = 2\pi$ . Then we have

$$q'(x)>0 \quad ext{for } x\in(0,\pi), \qquad q'(x)<0 \quad ext{for } x\in(\pi,2\pi), \ q'(0)=q'(\pi)=q'(2\pi)=0, \qquad q''(0)=q''(2\pi)=1, \qquad q''(\pi)=-1.$$

Moreover, since q is symmetric with respect to  $\pi$ , we can easily conclude that the thresholds corresponding to each subinterval has the same value, i.e.,  $\lambda_0 = \lambda_{01} = \lambda_{02}$ . Thus, according to Main Result, the threshold  $\lambda_0$  satisfies the inequalities

$$0.25 = rac{1}{4} < \lambda_0^2 \leq rac{243}{4} = 60.75.$$

Let us see how to improve the above-mentioned estimate by constructing a specific upper and lower functions.

Czech-Georgian Workshop Brno, July 2-4, 2024

15/17

Consider  $q(x) = 2 - \cos x$  for  $x \in [0, 2\pi]$ . Then local extremes of q divide the interval  $[0, 2\pi]$  into two subintervals, in particular, we set  $T = 2\pi$ ,  $x_1 = 0$ ,  $x_2 = \pi$ ,  $x_1 + T = 2\pi$ . Then we have

$$q'(x)>0 \quad ext{for } x\in(0,\pi), \qquad q'(x)<0 \quad ext{for } x\in(\pi,2\pi), \ q'(0)=q'(\pi)=q'(2\pi)=0, \qquad q''(0)=q''(2\pi)=1, \qquad q''(\pi)=-1.$$

Moreover, since q is symmetric with respect to  $\pi$ , we can easily conclude that the thresholds corresponding to each subinterval has the same value, i.e.,  $\lambda_0 = \lambda_{01} = \lambda_{02}$ . Thus, according to Main Result, the threshold  $\lambda_0$  satisfies the inequalities

$$0.25 = rac{1}{4} < \lambda_0^2 \leq rac{243}{4} = 60.75.$$

Let us see how to improve the above-mentioned estimate by constructing a specific upper and lower functions.

According to Corollary it is sufficient to find suitable functions  $\ell_1(x)$  and  $\ell_2(x)$  that satisfy (25)–(28). Obviously, we can start with positive constant functions.

Czech-Georgian Workshop Brno, July 2-4, 2024

15/17

Consider  $q(x) = 2 - \cos x$  for  $x \in [0, 2\pi]$ . Then local extremes of q divide the interval  $[0, 2\pi]$  into two subintervals, in particular, we set  $T = 2\pi$ ,  $x_1 = 0$ ,  $x_2 = \pi$ ,  $x_1 + T = 2\pi$ . Then we have

$$q'(x)>0 \quad ext{for } x\in(0,\pi), \qquad q'(x)<0 \quad ext{for } x\in(\pi,2\pi), \ q'(0)=q'(\pi)=q'(2\pi)=0, \qquad q''(0)=q''(2\pi)=1, \qquad q''(\pi)=-1.$$

Moreover, since q is symmetric with respect to  $\pi$ , we can easily conclude that the thresholds corresponding to each subinterval has the same value, i.e.,  $\lambda_0 = \lambda_{01} = \lambda_{02}$ . Thus, according to Main Result, the threshold  $\lambda_0$  satisfies the inequalities

$$0.25 = rac{1}{4} < \lambda_0^2 \leq rac{243}{4} = 60.75.$$

Let us see how to improve the above-mentioned estimate by constructing a specific upper and lower functions.

According to Corollary it is sufficient to find suitable functions  $\ell_1(x)$  and  $\ell_2(x)$  that satisfy (25)–(28). Obviously, we can start with positive constant functions. Then, if we put

$$egin{aligned} &\lambda_1^2 \stackrel{def}{=} \min\left\{(\ell_1^2 q''(x) + \ell_1 q(x))q^3(x): x \in [0,\pi]
ight\}, \ &\lambda_2^2 \stackrel{def}{=} \max\left\{(\ell_2^2 q''(x) + \ell_2 q(x))q^3(x): x \in [0,\pi]
ight\}, \end{aligned}$$

we can easily verify that the inequalities (26) and (27) with a = 0,  $b = \pi$  are fulfilled. Consequently, if also (28) is fulfilled, then we can conclude that (24) holds.

R. Hakl

As a first approximation we can put

$$\boldsymbol{\ell}_1 = \boldsymbol{\ell}_2 = \frac{3}{2}$$

because 3/2  $\in$   $(y_1(\lambda), y_2(\lambda))$ , no matter what  $\lambda^2 <$  60.75 is.

Czech-Georgian Workshop Brno, July 2-4, 2024

16/17

As a first approximation we can put

$$\ell_1=\ell_2=\frac{3}{2},$$

because  $3/2 \in (y_1(\lambda), y_2(\lambda))$ , no matter what  $\lambda^2 < 60.75$  is. Then, we get

$$\lambda_1^2 = 3.75, \qquad \lambda_2^2 = 60.75.$$

However, in this case  $\ell_2 = y_1(\lambda_2) = y_2(\lambda_2)$ , and so only the lower estimate is improved:

 $3.75 \leq \lambda_0^2 \leq 60.75.$ 

As a first approximation we can put

$$\ell_1=\ell_2=\frac{3}{2},$$

because  $3/2 \in (y_1(\lambda), y_2(\lambda))$ , no matter what  $\lambda^2 <$  60.75 is. Then, we get

$$\lambda_1^2 = 3.75, \qquad \lambda_2^2 = 60.75.$$

However, in this case  $\ell_2 = y_1(\lambda_2) = y_2(\lambda_2)$ , and so only the lower estimate is improved:

$$3.75 \le \lambda_0^2 \le 60.75.$$

Analyzing the function  $x \mapsto (\ell^2 q''(x) + \ell q(x))q^3(x)$  in more details, one can show that the optimal values for constant functions  $\ell_1$  and  $\ell_2$  are

$$\ell_1 = rac{20}{7}, \qquad \ell_2 = rac{5}{2}.$$

Czech-Georgian Workshop Brno, July 2-4, 2024

16 / 17

As a first approximation we can put

$$\ell_1=\ell_2=\frac{3}{2},$$

because  $3/2 \in (y_1(\lambda), y_2(\lambda))$ , no matter what  $\lambda^2 <$  60.75 is. Then, we get

$$\lambda_1^2 = 3.75, \qquad \lambda_2^2 = 60.75.$$

However, in this case  $\ell_2 = y_1(\lambda_2) = y_2(\lambda_2)$ , and so only the lower estimate is improved:

$$3.75 \leq \lambda_0^2 \leq 60.75.$$

Analyzing the function  $x \mapsto (\ell^2 q''(x) + \ell q(x))q^3(x)$  in more details, one can show that the optimal values for constant functions  $\ell_1$  and  $\ell_2$  are

$$\ell_1 = rac{20}{7}, \qquad \ell_2 = rac{5}{2}.$$

Then, we get

)

$$\lambda_1^2 = rac{540}{49} pprox 11.020408163, \qquad \lambda_2^2 = rac{3125}{64} = 48.828125,$$

 $y_1(\lambda_2)pprox 0.835507015894,$ 

Czech-Georgian Workshop Brno, July 2-4, 2024

and we have the estimate

$$\frac{540}{49} \le \lambda_0^2 \le \frac{3125}{64}.$$

R. Hakl

Periodic Travelling Waves

16 / 17

Let us pass to nonconstant functions  $\ell_1(x)$  and  $\ell_2(x)$ . Then the choice

$$\ell_i(x) \stackrel{def}{=} a_i + b_i \cos x + c_i \sin x + d_i \sin x \cos x \qquad ext{for } x \in [0,\pi] \quad (i=1,2),$$

where

$a_1 = 4.265,$	$b_1 = 1.639,$	$c_1 = -1.075,$	$d_1 = -0.778,$
$a_2 = 3.605,$	$b_2 = 1.025,$	$c_2 = -0.408,$	$d_2 = -0.222,$

guarantees that  $\ell_1(x)$  and  $\ell_2(x)$  satisfy (25)–(28) with  $a = 0, b = \pi, \lambda_1^2 = 26.4$ , and  $\lambda_2^2 = 31.68$ .



Let us pass to nonconstant functions  $\ell_1(x)$  and  $\ell_2(x)$ . Then the choice

$$\ell_i(x) \stackrel{def}{=} a_i + b_i \cos x + c_i \sin x + d_i \sin x \cos x \qquad ext{for } x \in [0,\pi] \quad (i=1,2),$$

where

$a_1 = 4.265,$	$b_1 = 1.639,$	$c_1 = -1.075,$	$d_1 = -0.778,$
$a_2 = 3.605,$	$b_2 = 1.025,$	$c_2 = -0.408,$	$d_2 = -0.222,$

guarantees that  $\ell_1(x)$  and  $\ell_2(x)$  satisfy (25)–(28) with a = 0,  $b = \pi$ ,  $\lambda_1^2 = 26.4$ , and  $\lambda_2^2 = 31.68$ . Furthermore, note also that

$$\ell_1(\pi) < y_2(\lambda_1), \qquad y_2(\lambda_2) < \ell_2(\pi)$$
 (29)

where  $y_2(\lambda_i)$  is the greater root of (19) with  $\lambda = \lambda_i$  (i = 1, 2). Indeed,

 $\ell_1(\pi) = 2.626, \quad y_2(\lambda_1) \approx 2.62792828771, \quad y_2(\lambda_2) \approx 2.53762549442, \quad \ell_2(\pi) = 2.58$ 

The condition (29) is stronger than (28) and allows strict inequalities in the threshold estimate. Therefore, according to Corollary we have

$$26.4 < \lambda_0^2 < 31.68.$$

R. Hakl